A panoramic view of Edinburgh, Scotland, featuring the Temple of Antonine in the foreground. The temple is a circular structure with four columns and a pediment, topped with a statue. The cityscape in the background includes various buildings, a clock tower, and a bridge, all under a blue sky with scattered white clouds.

Edinburgh, Scotland

<http://www.waimhcongress.org/>

Modelling of Premixed Flame under Harmonic Oscillation

Donghyuk Shin

Edward S Richardson

University of Edinburgh

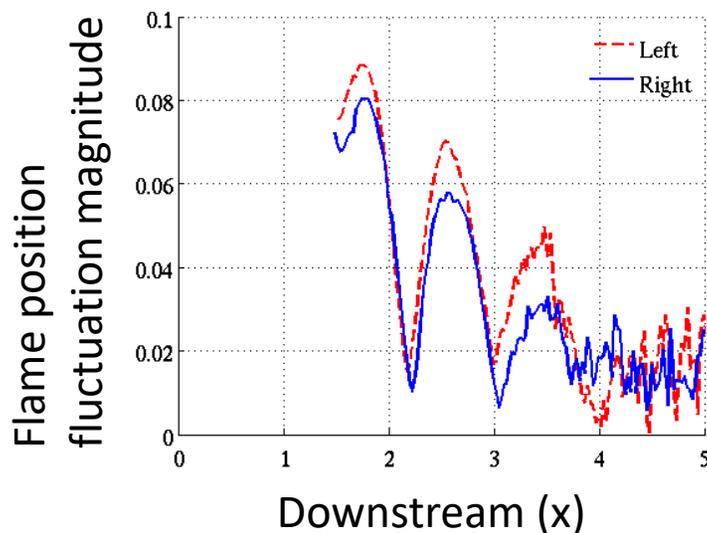
University of Southampton

13/9/2018, UKCTRF Meeting, Cambridge

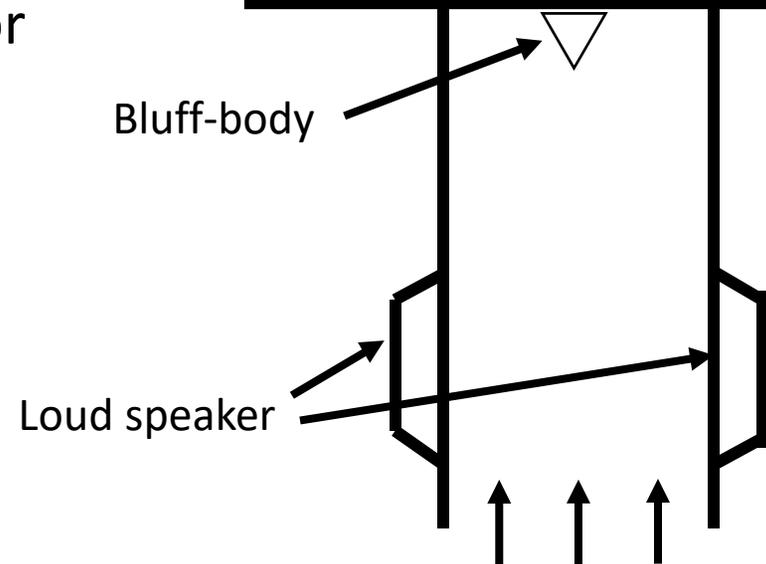
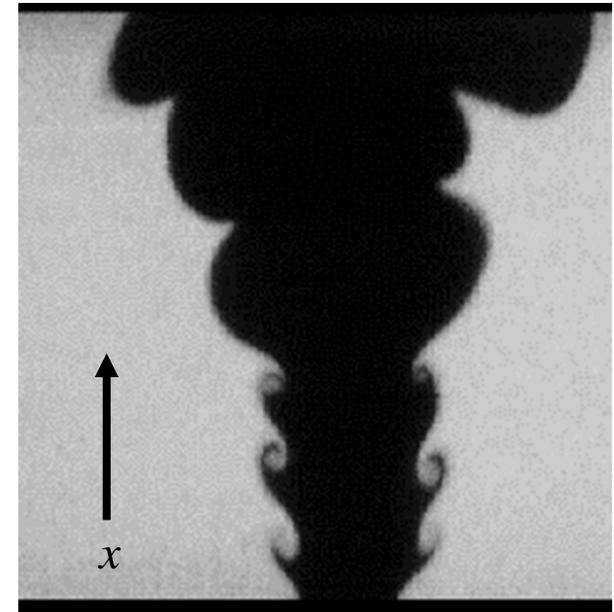
Introducing Flame-Holder Oscillation

Conventional ways to force flames

- Conventionally, we force the flow to study combustion instability
 - Developed flame wrinkle is not sinusoidal
 - Flow field continues to perturb flames, leading to interference behavior



Mie-scattering of flame



Bluff-body flame, velocity-focused, courtesy of Shanbhogue



Flame-holder oscillation

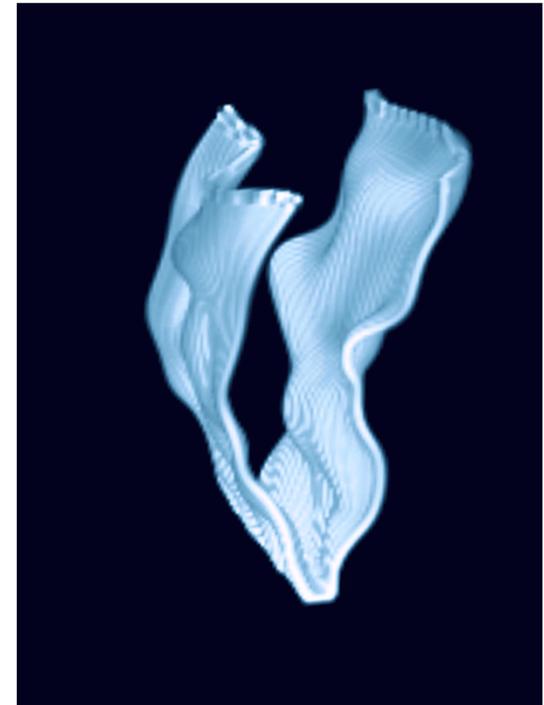
- Instead, we vibrate the flame holder
 - Flow field is nominally uniform
 - The flame wrinkle starts from a sine wave shape



Experimental setup
of Humphrey et al.



Mie-scattering of the flame



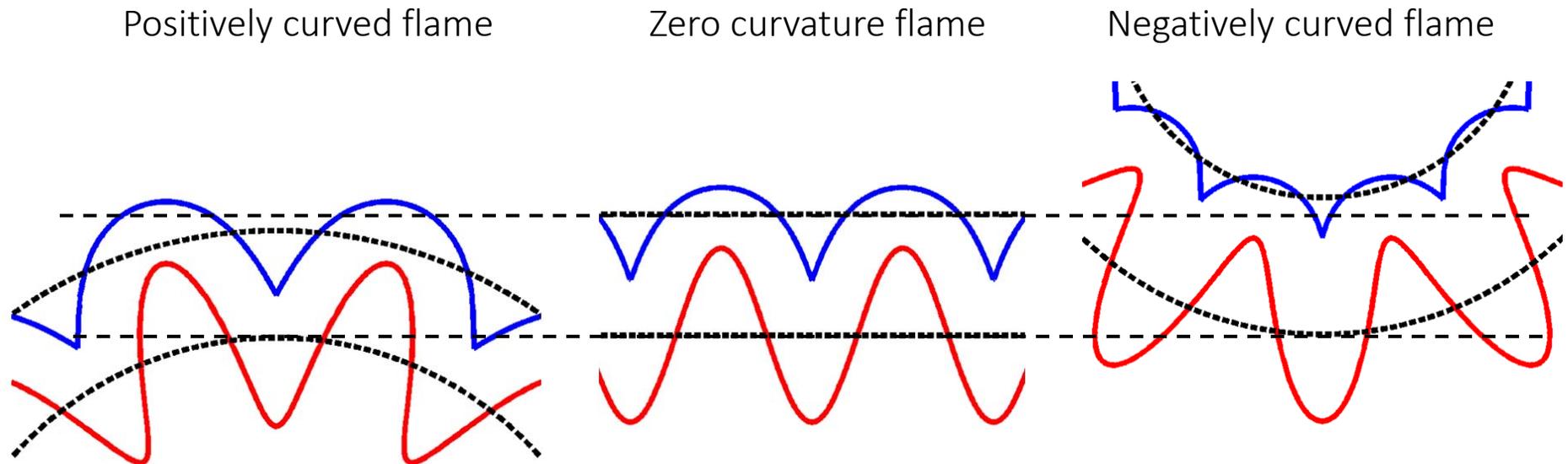
On-going DNS
By Shin, Aspden, Lieuwen



Turbulent Markstein Length

Turbulent Markstein Length: Curvature effect due to small scale wrinkles

- Consider a condition that laminar flame speed is constant, i.e., independent of curvature.
- Supposed 3 different initial flame positions



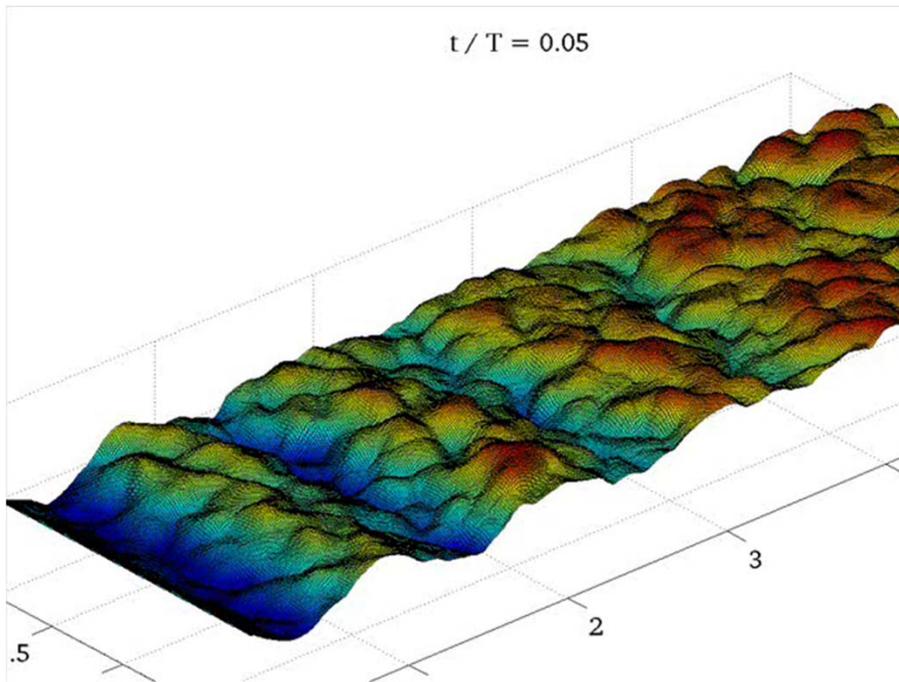
- Propagation speed, on average, depends on the curvature of ensemble averaged front



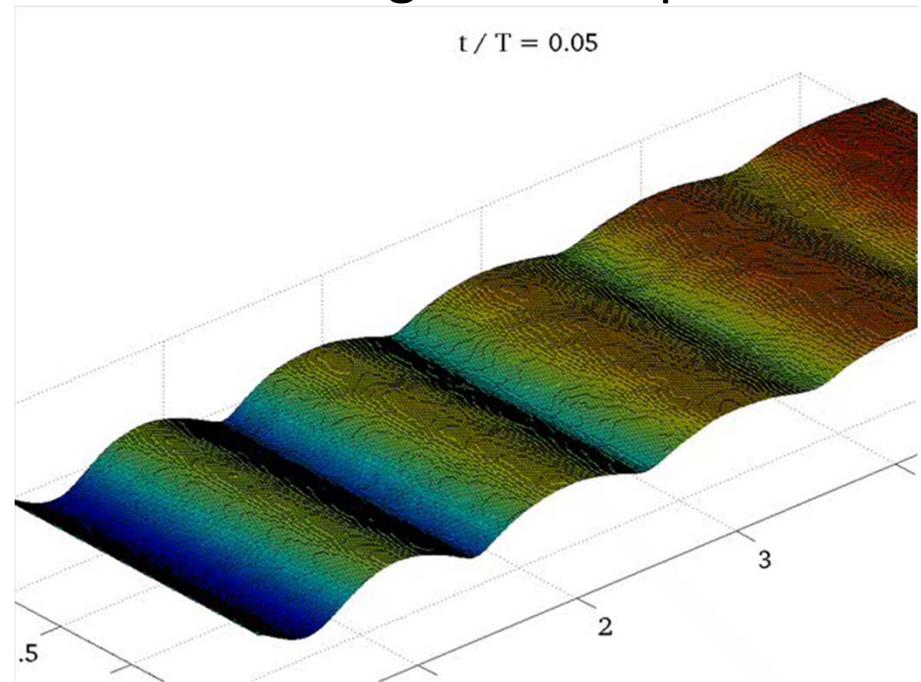
Turbulent Markstein length

- **Turbulent flame speed** depends on the **averaged curvature**

Instantaneous flame position



Phase-averaged flame position



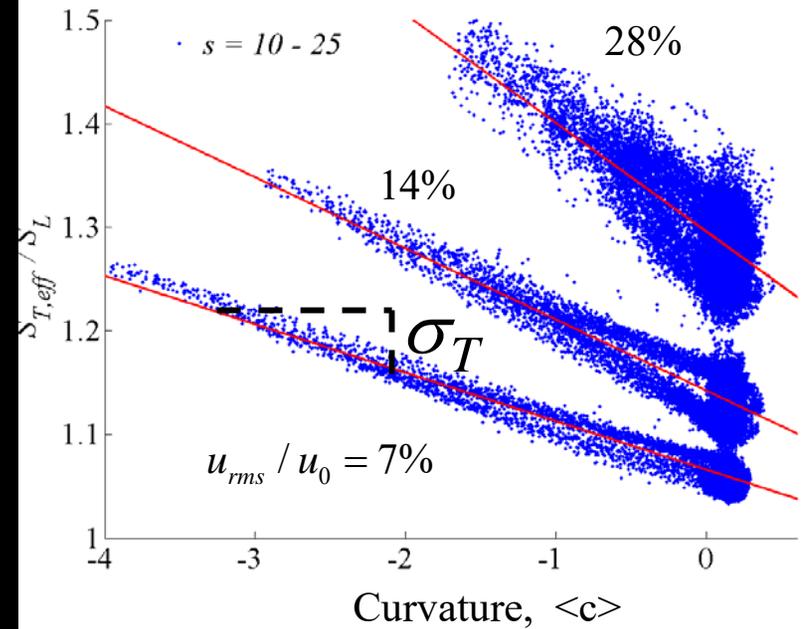
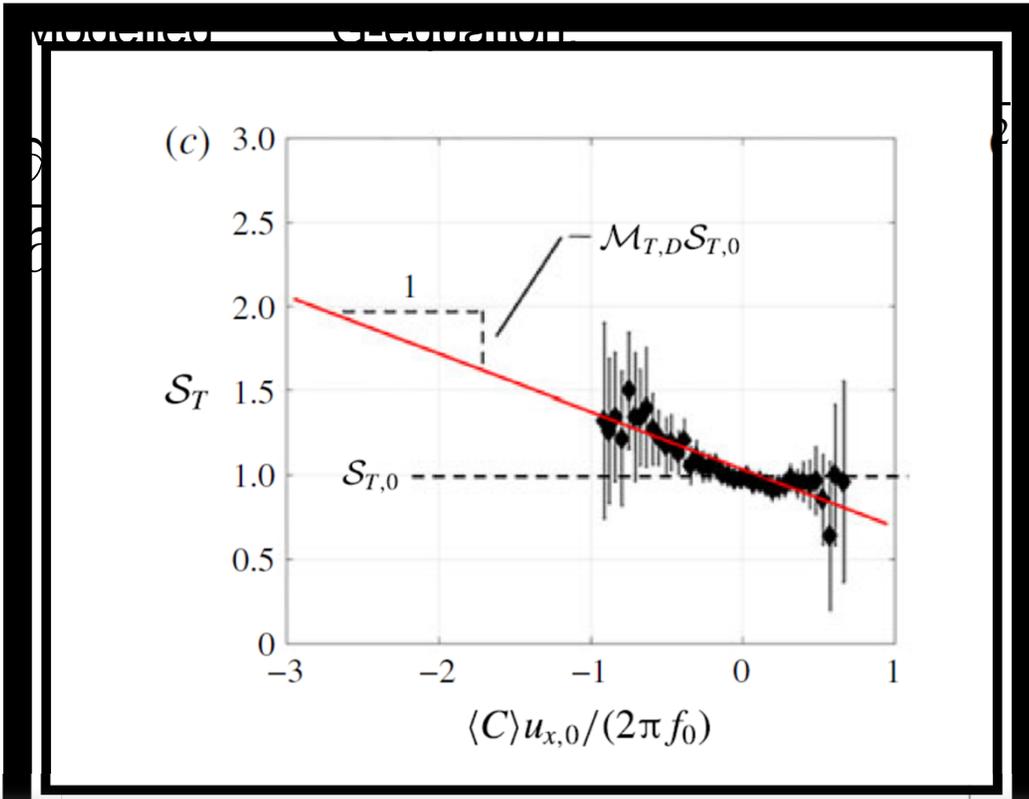
By 3D G-equation simulation



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$$u_{rms} / u_0 = 7 \%$$

Turbulent Flame Speed

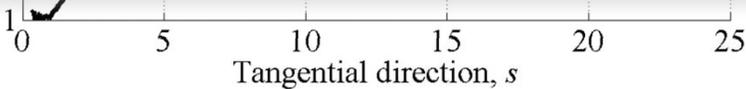


Nearly linear correlation with flame

This dependency is verified experimentally recently (Humphrey, Emerson, Lieuwen, JFM, 2018)

“Karlovitz Length,” σ_T

$$S_{T,eff} = S_{T,eff} (1 - \sigma_T \langle c \rangle)$$



Quantifying
the Curvature effect

High order asymptotic approach

Example:
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0$$

- First order asymptotic method:

If we decompose the velocity into: $u = u_0 + u'$

$$\frac{\partial u'}{\partial t} + u_0 \frac{\partial u'}{\partial x} = 0$$

Which has an solution: $u' = u'_{init} \left(x - \frac{t}{u_0} \right)$

- Second order asymptotic method:

Expand in the orders of fluctuation magnitudes: $u = u_0 + \varepsilon u_1 + \varepsilon^2 u_2$

$$\varepsilon \left[\frac{\partial u_1}{\partial t} + u_0 \frac{\partial u_1}{\partial x} \right] + \varepsilon^2 \left[\frac{\partial u_2}{\partial t} + u_1 \frac{\partial u_1}{\partial x} + u_0 \frac{\partial u_2}{\partial x} \right] + O(\varepsilon^3) = 0$$

$$O(\varepsilon): \quad \frac{\partial u_1}{\partial t} + u_0 \frac{\partial u_1}{\partial x} = 0$$

$$O(\varepsilon^2): \quad \frac{\partial u_2}{\partial t} + u_0 \frac{\partial u_2}{\partial x} = -u_1 \frac{\partial u_1}{\partial x}$$

Get the solution for u_1 .

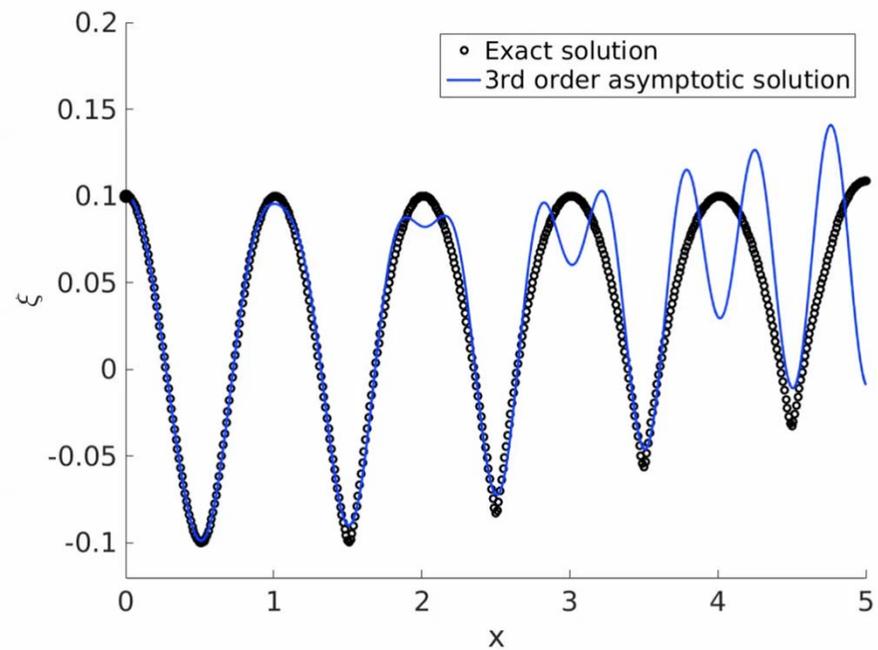
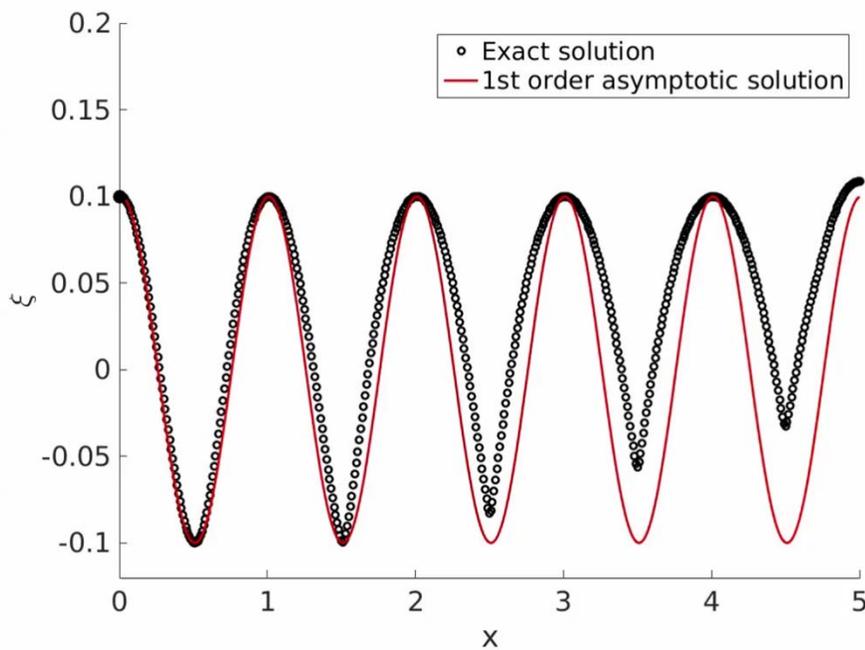
Get the solution for u_2 .



Comparison of exact vs asymptotic solutions

1st order asymptotic solution

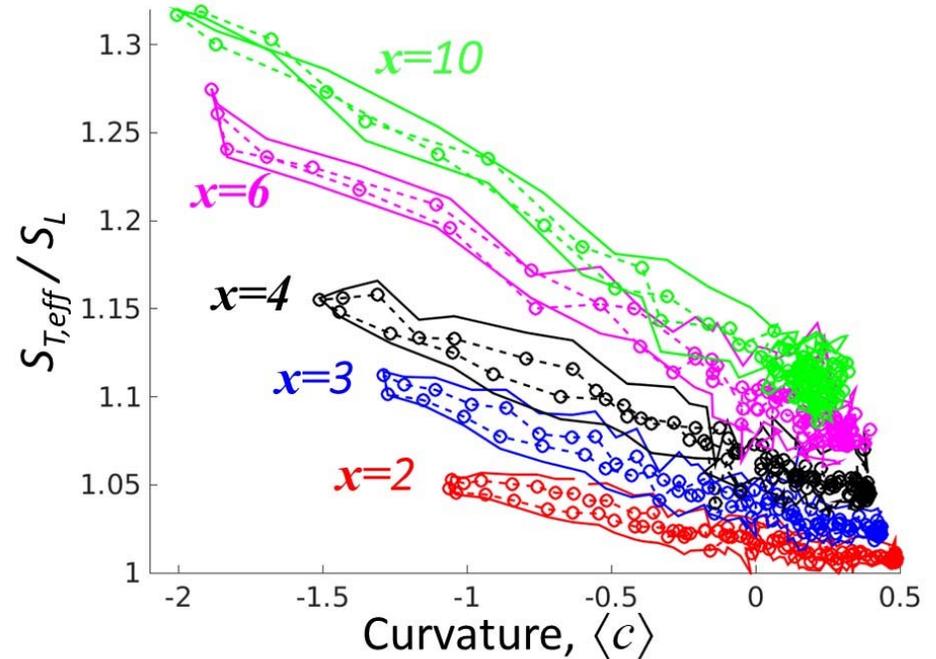
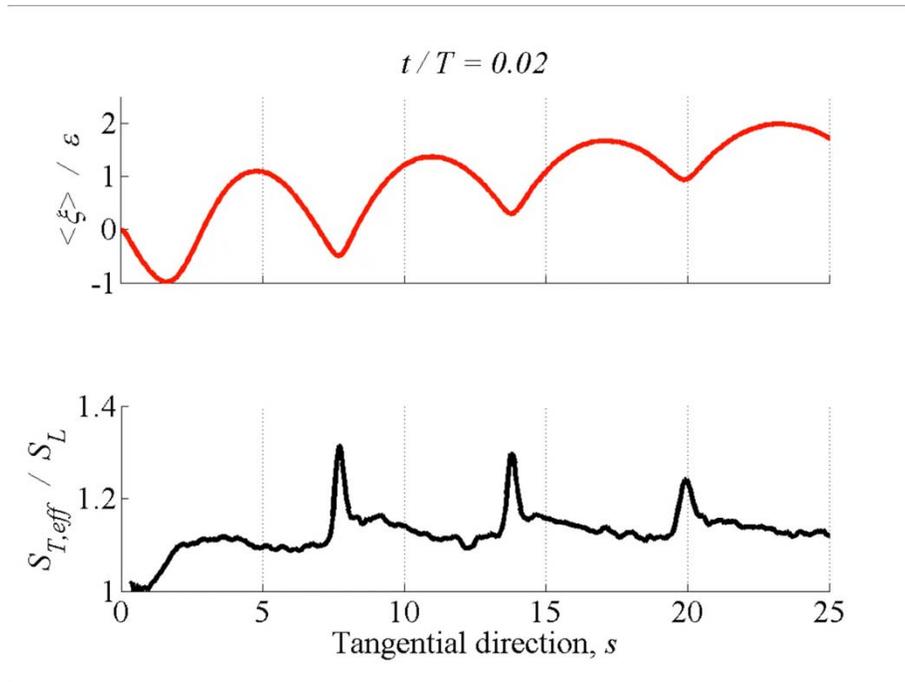
3rd order asymptotic solution



Higher order asymptotic solution works better near $x=0$, but deviates more further downstream.



Turbulent Markstein length at different locations



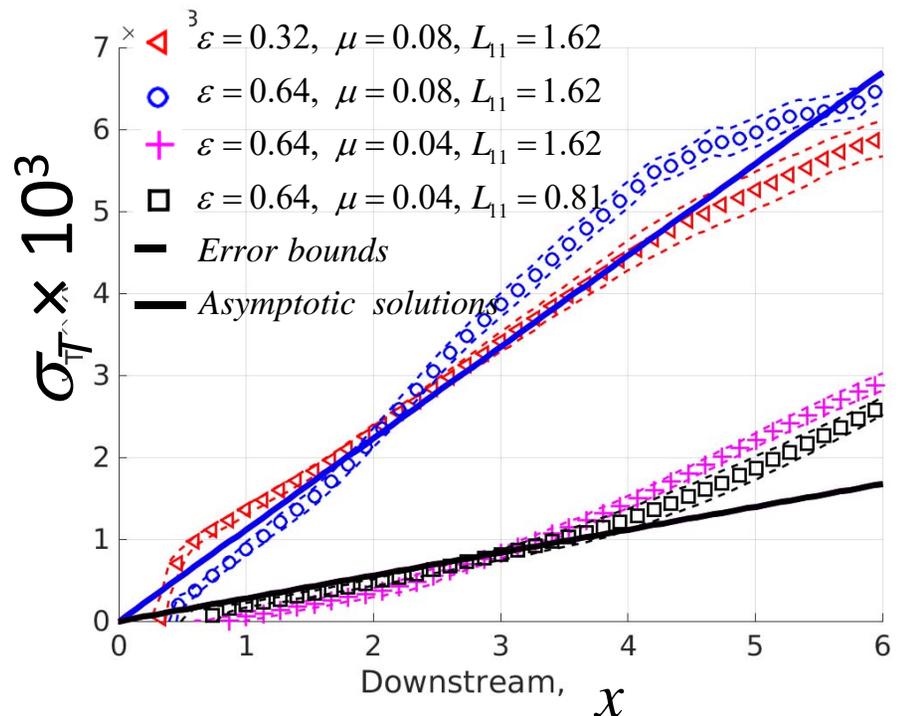
Turbulent Markstein Length for Iso-thermal condition

- Simulations done by solving G-equation with frozen flow
- Asymptotic expansion up to the 3rd order

$$\sigma_T \approx \mu^2 \left(\frac{S_L}{u_0} + \frac{u_0}{S_L} \right) x$$

μ : turbulent intensity

where $S_{T,eff} = S_{T,0} (1 - \sigma_T \langle C \rangle)$



Far field asymptotic solution for Turbulent Markstein length

$$\sigma_T = \left[\left\langle \frac{\partial \xi'}{\partial t} \frac{\partial v'}{\partial y} \right\rangle - \frac{1}{2} \frac{\partial \langle v'^2 \rangle}{\partial y} \right] - \frac{1}{S_L} \left[\left\langle \xi' \frac{\partial^2 v'}{\partial y^2} \right\rangle + \left\langle \frac{\partial \xi'}{\partial t} \frac{\partial u'}{\partial y} \right\rangle - \left\langle \frac{\partial \xi'}{\partial z} \frac{\partial w'}{\partial y} \right\rangle - \frac{\partial \langle u'v' \rangle}{\partial y} \right]$$

u', v', w' : fluctuating velocity components
 ξ' : fluctuating flame position

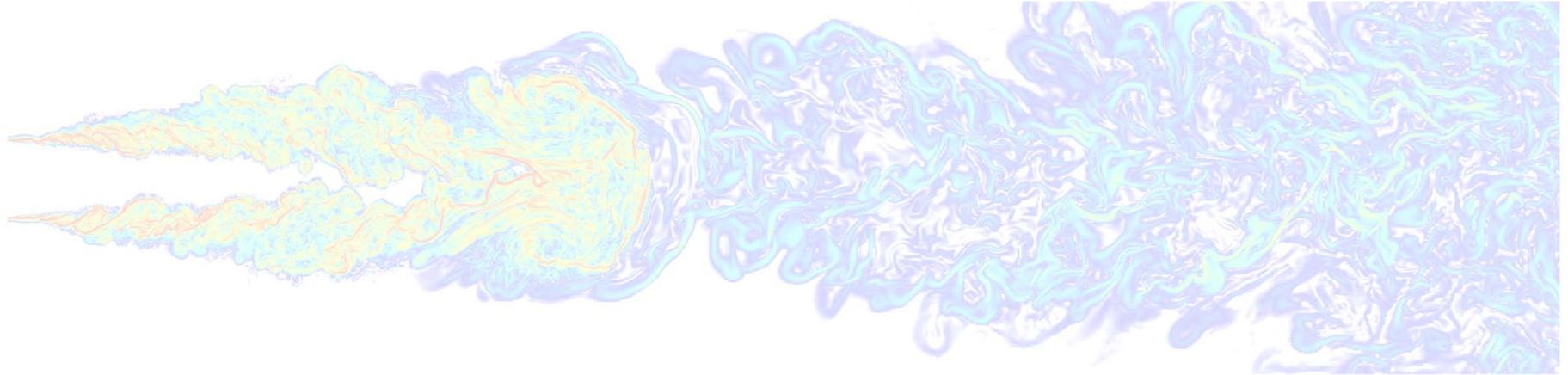
- This has not been validated with simulations yet.



Conclusion

- Oscillating flame holder setup is an interesting case to study unsteady flame dynamics
- Turbulent flame speed can depends on curvature
- Asymptotic analysis indicates that the turbulent Markstein length is proportional to squared turbulent intensity and laminar flame speed





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