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Flame-Flame Interaction and Thermoacoustic Instability

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Motivation

- Flame-Flame interaction can :
 - Induce unstable mode bifurcation (J. Dawson et al. 2014)
 - Increase heat release (B. Ranganath et. al. 2006)
 - Change the symmetry of a system (N. Worth et al. 2012)
 - Create a major source of noise (M. Talei et al. 2012)
 - Influence driving forces for thermoacoustic instability (R. Balachandran et al. 2005)
 - Induce local changes in topology Today's focus

•How do the flow/scalar topologies compare between a single flame and duel flames?

•Can multiple flames be modelled with a single equivalent?

Type of interaction

- Normal Interaction:
 - -Initial acceleration
 - -Merging
 - -Cusp formation (negative curvature)
 - -Cusp recovery (loss of flame surface)
- Counter-normal Interaction:
 - -Against normal direction of propagation-Driven by turbulence/high positive curvature-Longer time scale



Source: T. Dunstan, N. Swaminathan et al. 2013

Flow and Scalar Topologies



L. Cifuentes et al. 2016

W. Han et al 2019

Numerical Methodology

- CompReal- inhouse, finite difference code
- Fully compressible, density based solver
- Skew-symmetric (4th order) discretisation
- High order Runge-Kutta integration in time
- Chemkin interface
- Navier-Stokes Characteristic Boundary Conditions
- Turbulent inflow generator
- LES with modified Artificially Thickened Flame (ATF) closure need to preserve gradients in regions of flame

ATF

"Classical"

• Transformation: $\xi(x) = \int \mathcal{F}(x') dx'$

$$\frac{\partial \rho Y}{\partial \tau} + \frac{\partial \rho u_i}{\partial \xi_i} - \frac{\partial}{\partial \xi_i} (\rho D \mathcal{F} \frac{\partial Y}{\partial \xi_i}) - \frac{\dot{\omega}}{\mathcal{F}} = 0$$

- •Reduced effective source term "thickens" the flame but maintains correct flame speed
- •Alters nature of turbulence/flame interaction
- •Efficiency factor needed to model sub grid flame wrinkling

"Modified"

•Thicken momentum and energy transport equations as well – maintains correct Damkohler number

$$\frac{\partial \bar{\rho} \tilde{u}_i}{\partial \tau} + \frac{\partial \bar{\rho} \tilde{u}_i \tilde{u}_j}{\partial \xi_j} + \frac{\partial \bar{P}}{\partial \xi_i} = \frac{\partial 2\mathcal{F}\mu \hat{S}^*_{ij}}{\partial \xi_j} + (1 - \Omega) \frac{\partial \tau^{sgs}_{ij}}{\partial \xi_j}$$
$$\frac{\partial \bar{\rho} \tilde{Y}_k}{\partial \tau} + \frac{\partial \bar{\rho} \tilde{u}_j \tilde{Y}_k}{\partial \xi_j} = \frac{\partial}{\partial x_j} \left(\bar{\rho} \mathcal{F} D \frac{\partial \tilde{Y}_k}{\partial \xi_j} \right) + (1 - \Omega) \frac{\partial J^{sgs}_j}{\partial \xi_j} + \frac{\dot{\omega}}{\mathcal{F}}$$
$$\mathcal{F} = 1 + (\mathcal{F}_0 - 1)\Omega$$

•No need for efficiency factor

Numerical Setup

- Premixed methane/air flames at unity equivalence ratio
- •30mm flame separation distance
- •Two different bulk velocities simulated (12 and 16 m/s) (add re and regime)
- •A single flame at 12m/s also simulated for comparison
- •6 million mesh points per flame
- •Used grid stretching to attain resolution of the order of unthickened laminar flame thickness
- •After thickening have at least 10 points within flame



Mean Curvature

- •Skewness factor of -0.57 compared with -0.52 from experiment
- •Sampled over at least 10 eddy turnover times
- Rescaled into physical dimensions via thickening factor
- •Rescaling does not affect *type* of scalar toplogy



Tyagi, Ankit, et al. Combustion and Flame (2019).

Scalar Joint PDFs



Scalar Structure

Most probable topology is "nearly flat" as expected

•As velocity increases the distribution becomes wider (i.e. more wrinkling)

•Single flame favours slightly convex structures over slightly concave structures – this bias is decreased in duel flame cases

•Experiments suggest single flame has higher degree of normal flame-flame interaction than the duel case and vice-versa for counter-normal

•Overall, across all scales concave structures are more likely

Flow Topologies

•S2 most dominant as expected

•S4 and S7 disappear due to destructive of vorticity by positive dilatation

 However S1 and S5 increase when going from reactant to product side – especially for duel flame case. Most likely this is due to the "opening" of the flames increasing compressive topologies

•When adding a second flame, note the increase in S2 and decrease in S3



Influence of flow on scalar topologies

•S2- Mostly concave

- •S3- Mostly convex
- •Decrease in S3 in duel flames shown earlier could be one reason for preference of concave structures
- •Proportion of concave structures is 51%,58% and 60% when compared across all scales for the single, duel slow and duel fast flames respectively



Impact of dilatation

•So why does S3 reduce in duel flames?

•S3 occurs in regions of very small dilatation

- •Adding a second flame increases mean dilatation
- •This could be one of the driving mechanisms for the reduction of S3 observed for duel flames



Conclusions

•Overall all flames prefer concave scalar structures over convex

•This agrees well with the experimental findings of normal flame-flame interactions being more prevalent than counter-normal

•Structures which were "nearly" flat showed the opposite trend and favoured convex structures locally for a single flame, this requires further investigation

•Compressive flow topologies not commonly observed in premixed flame were seen to persist across the flame – most likely due to the "opening" effect

•The main effect of adding a second flame was a reduction in convex structures

•This reduction can be attributed at least partly due to the changes in flow topology induced by dilatation

Future work

•Currently working on using AMR to investigate the effect of varying thickening factor

•Effects of varying flame spacing/acoustic forcing

•Extension to multiple flames and application to global stability





Any Questions?

THANK YOU FOR YOUR TIME